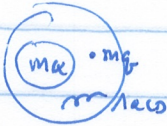


31 MAY 2016 (am)

SCE1 in QCD

Lecture 2: jets & building blocks of SCE1

... Kinish HQET:



$$p_0 = m_q v + k$$

$v \cdot k_{QCD} \ll m_q$

$$\Rightarrow i \frac{1+\not{v}}{2} \frac{1}{v \cdot k + i\epsilon} \xrightarrow{\text{split}} \frac{\not{\epsilon} k}{\epsilon} \text{ig } U_\mu^T A$$

construct Lag: step 1:  $\psi(x) = e^{-im_q v \cdot x} \psi_v(x)$

step 2:  $\psi_v = h_v + H_v$

$$h_v = \frac{1+\not{v}}{2} \psi_v \quad H_v = \frac{1-\not{v}}{2} \psi_v$$

step 3:  $\mathcal{L}_{\text{Dirac QCD}} = \bar{\psi} (i\not{D} - m_q) \psi$

$$= \bar{\psi}_v e^{im_q v \cdot x} (i\not{D} - m_q) e^{-im_q v \cdot x} \psi_v$$

$$= \bar{\psi}_v (i\not{D} + m_q \not{v} - m_q) \psi_v$$

$$(1-\not{v}) h_v = 0$$

$$(1+\not{v}) H_v = 2 H_v$$

$$= (\bar{h}_v + \bar{H}_v) [i\not{D} h_v + (i\not{D} - 2m_q) H_v]$$

$$\frac{1+\not{v}}{2} i\not{D} \frac{1+\not{v}}{2} = \frac{1+\not{v}}{2} i\not{v} \cdot D \quad \frac{1-\not{v}}{2} i\not{D} \frac{1-\not{v}}{2} = \frac{1-\not{v}}{2} (-i\not{v} \cdot D)$$

$$= \bar{h}_v i\not{v} \cdot D h_v - \bar{H}_v (i\not{v} \cdot D + 2m_q) H_v$$

$$+ \bar{h}_v i\not{D} H_v + \bar{H}_v i\not{D} h_v$$

$$(\bar{h}_v m_q H_v = 0)$$

actually only  $\not{D}_\perp$  where  $X_\perp^\mu = X^\mu - X \cdot v v^\mu$

step 4:  $H_v$  is "heavy", mass  $2m_q$

$\Rightarrow$  integrate out, solve classical eq. of motion

$$\Rightarrow (i\not{v} \cdot D + 2m_q) H_v = i\not{D}_\perp h_v$$

$$\Rightarrow \mathcal{L} = \bar{h}_v (i\not{v} \cdot D + i\not{D}_\perp \frac{1}{i\not{v} \cdot D + 2m_q} i\not{D}_\perp) h_v$$

$$= \bar{h}_v (i\not{v} \cdot D + \frac{1}{2m_q} \not{D}_\perp \not{D}_\perp + \dots) h_v$$

$$= \bar{h}_v (i\not{v} \cdot D - \frac{\not{D}_\perp^2}{2m_q} - g \frac{\sigma^{\mu\nu} G_{\mu\nu}}{2m_q}) h_v + \dots$$

$\Rightarrow$  reproduces  $\frac{g}{2m_q}$  Fermi rules  $\leftarrow \begin{matrix} \mathcal{O}(1) & \mathcal{O}(1/m_q) & \mathcal{O}(1/m_q^2) \end{matrix}$

$$\begin{aligned} & \frac{1+\not{v}}{2} i\not{D} \frac{1+\not{v}}{2} \\ &= \frac{1+\not{v}}{2} (i\not{v} \cdot D - \not{v} \cdot D) \\ &= \frac{1+\not{v}}{2} i\not{v} \cdot D \\ &= \frac{1+\not{v}}{2} i\not{D}_\perp \end{aligned}$$

$$\begin{aligned} & \frac{1-\not{v}}{2} i\not{D} \frac{1-\not{v}}{2} \\ &= \frac{1-\not{v}}{2} (i\not{v} \cdot D + \not{v} \cdot D) \\ &= \frac{1-\not{v}}{2} (-i\not{v} \cdot D) \\ &= -\frac{1-\not{v}}{2} i\not{D}_\perp \end{aligned}$$

(some higher order terms constrained by RPI)

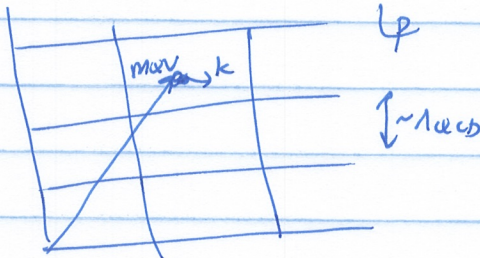
$$p_a = m_a v^i k^j \quad \text{inv. to} \quad v \rightarrow v + \frac{e}{m_a} k$$

$$k \rightarrow k - e$$

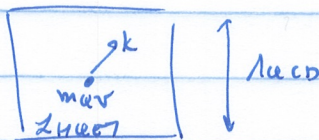
if more than one heavy quark e.g.  $(B \rightarrow D \bar{D}^*)$   $\frac{b}{v} \frac{c}{v'}$

picture in momentum space:

$$\mathcal{L}_{\text{HQET}} = \sum_v \bar{\psi}_v \not{D} \psi_v$$

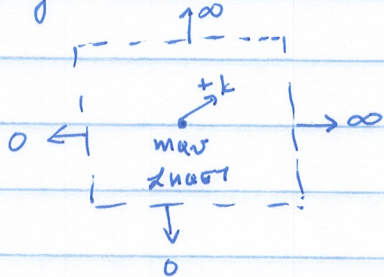


HQET lives in the box



but cutoff violates Lorentz & gauge invariance

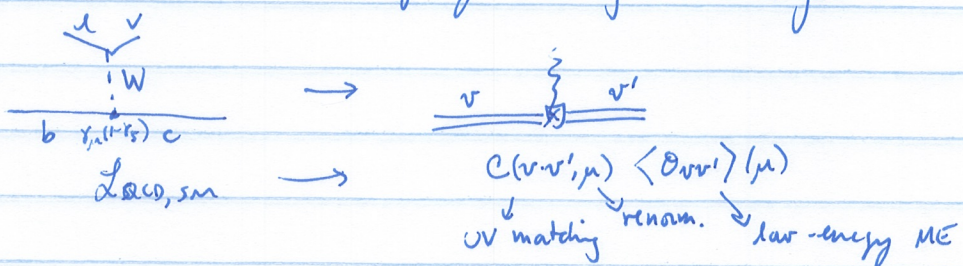
we use dim. reg & scale boundaries away.



$\mathcal{L}_{\text{HQET}}$  reproduces QCD in the box for  $k \sim \Lambda_{\text{QCD}}$  (IR)

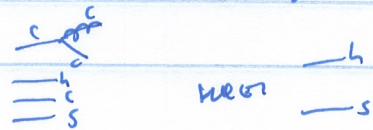
but is bad for large  $k \rightarrow \infty$  and different UV divergence structure as  $k \rightarrow \infty$

↳ take care of by matching & renormalization



in SCET will find

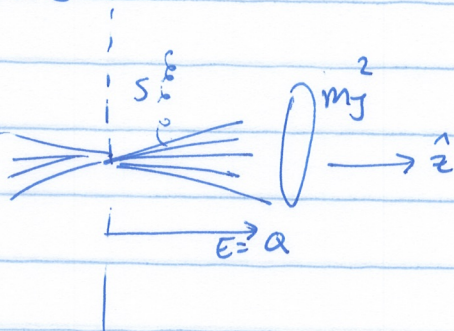
- labels can change
- multiple scales



## II Jets & SCET bldg. blocks

Degrees of freedom:

measure jet mass



light-cone coords:

$$n = (1, \hat{z}) \quad n^2 = \bar{n}^2 = 0$$

$$\bar{n} = (1, -\hat{z}) \quad n \cdot \bar{n} = 2$$

$$P^\mu = \frac{\bar{n} \cdot p}{2} n^\mu + \frac{n \cdot p}{2} \bar{n}^\mu + P_L^\mu$$

$$\approx P_L^\mu + \dots$$

$$P^2 = \bar{n} \cdot p \, n \cdot p + P_L^2$$

$$= p^+ p^- - \vec{P}_L^2$$

$$P = (p^-, p^+, P_L)$$

$$P^- = \bar{n} \cdot p = E + p_z$$

$$P^+ = n \cdot p = E - p_z$$

measure inv. mass of particles in a hemisphere:

$$P_c^2 = p^+ p^- + P_L^2 \sim m^2$$

$$\Rightarrow p^+ \sim \frac{m^2}{Q}$$

$$P_c \sim (Q, \frac{m^2}{Q}, m)$$

$$\approx Q(1, \lambda^2, \lambda)$$

but I can add:

$$P_c + P_s$$

$$\lambda \sim m/Q$$

$$P_s \sim (\frac{m^2}{Q}, \frac{m^2}{Q}, \frac{m^2}{Q}) = Q(\lambda^2, \lambda^2, \lambda^2)$$

and preserve jet mass

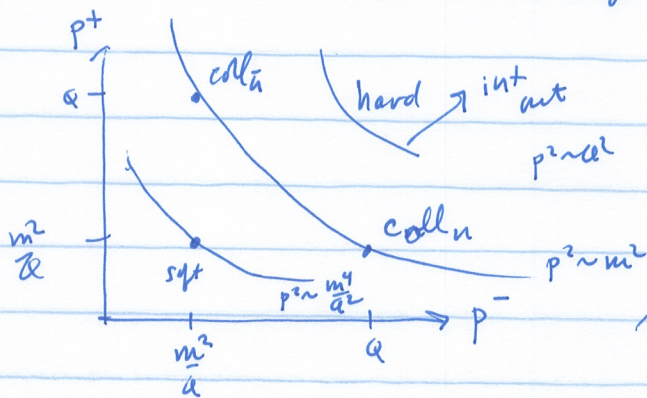
$$\frac{P_s}{P_c} \sim \frac{m^2}{m^2} \sim 1$$

$$P_J \sim (Q, \frac{m^2}{Q}, \frac{m^2}{Q})$$

c,s "talk" through  $p^+$  comp. only

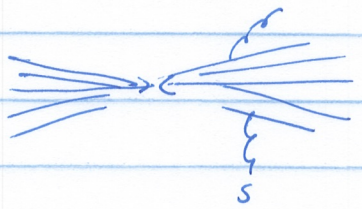
modes in "SCET<sub>I</sub>":

e.g. back-to-back jets



separated in virtuality  
dim reg & scale  $\mu$  alone regulates.

consider measuring jet / beam  $P_T$ :



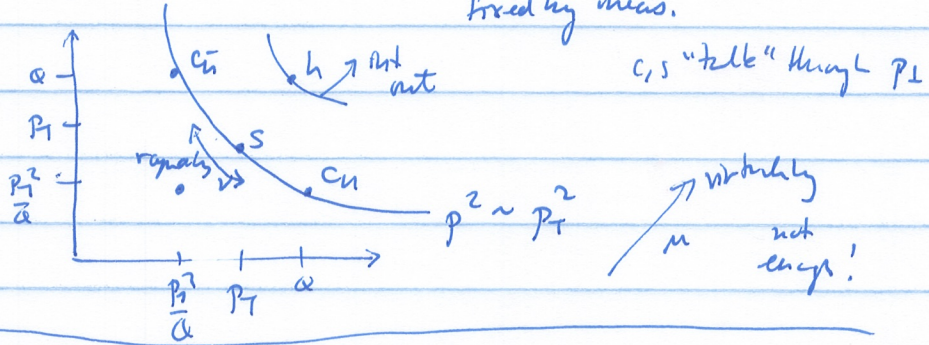
$$P_c = (\alpha, \frac{P_T^2}{\alpha}, P_T)$$

$$= \alpha(1, \eta^2, \eta) \quad \eta \sim \frac{P_T}{Q}$$

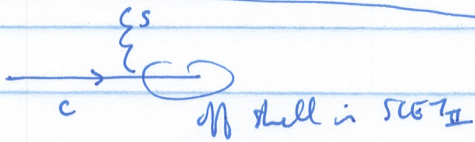
$$P_s = (P_T, P_T, P_T) = \alpha(\eta, \eta, \eta)$$

↑ fixed by meas.

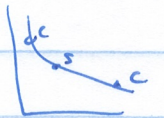
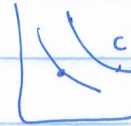
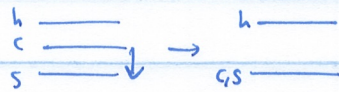
modes in "SCET<sub>II</sub>",



end lecture 2



can match  $QCD \rightarrow SCET_I \rightarrow SCET_{II}$



lat collinear momenta:

$$P_c = \tilde{P}_c + k$$

↓ label
↓ residual  $\sim \alpha \lambda^2$

Use MeV in HQET

but now 2 pieces:

$$\tilde{P}_c = \bar{n} \cdot \tilde{P}_c \frac{n}{2} + \tilde{P}_\perp$$

$\downarrow \sim \alpha$ 
 $\downarrow \sim \alpha \lambda$

collinear fields:

$$\psi_c \rightarrow \psi_c(x) = \sum_{\tilde{p} \neq 0} e^{-i\tilde{p} \cdot x} \psi_{n,p}(x)$$

↳ sum (not in HQET) needed because of coll. gluons

$$A^\mu \rightarrow A_c^\mu = \sum_{\tilde{p} \neq 0} e^{-i\tilde{p} \cdot x} A_{n,p}^\mu(x)$$

$$+ A_s^\mu \quad \sim \alpha(\lambda^1, \lambda^1, \lambda^2)$$

label  $\tilde{p} \neq 0$   
avoid overlap with soft