

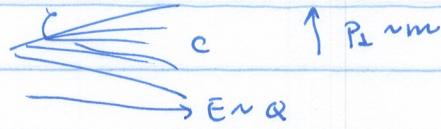
31 MAY 2016 (pm)

Lecture 3: Construction of SCET_I Lagrangian

- mimic steps in HQET:
- 1) factn out large mass. phase
 - 2) project out "large" & "small" components
 - 3) int.-act "heavy" field (e.g. of mott.)
 - 4) expand λ in powers of λ
 - 5) new to SCET: coll. Wilson lines,
"label operators"

coll & soft d.o.P's e.g. in jets

$$s \{ k \cdot v^2 / Q^2$$

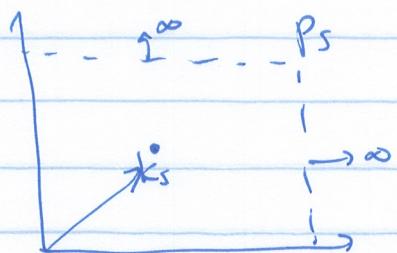
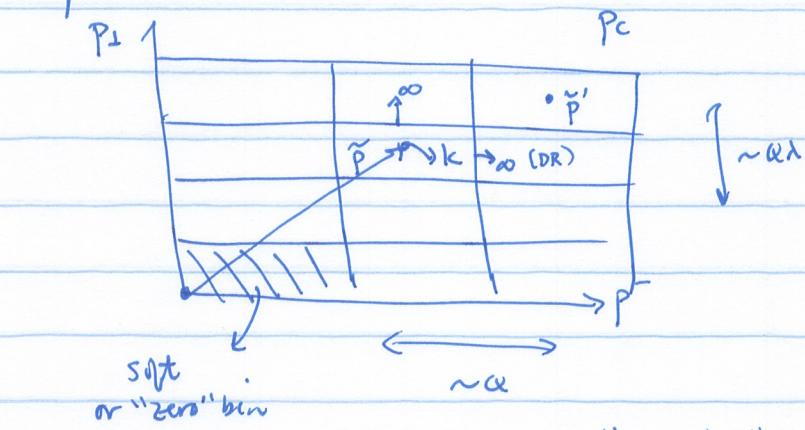


split $p_c = \tilde{p}_c + k$
 \downarrow
 residual $\sim Q\lambda^2$

label $\tilde{p} = \tilde{n} \cdot \tilde{p} \frac{n^\mu}{2} + \tilde{p}_\perp$
 \downarrow
 $\partial(\alpha)$ \downarrow $\partial(\alpha\lambda)$

like HQET mark
2 scales in label

momentum space will look like:



will need "new bin" subtraction
in coll. sector

odd amp.

$$\sum_{p \neq 0} \int d^4k \rightarrow \int_0^\infty d^4p (\delta_{p_c} - \delta_{p_0})$$

discrete columns (DR)

$p_c \rightarrow 0$ bin

$p_c \rightarrow 0$ bin

$$\text{coll. prop. } \frac{p_c}{\gamma} = \frac{i(\tilde{p} + k)}{(p+k)^2 + i\epsilon} = \frac{i\bar{n}\cdot\tilde{p}}{\bar{n}\cdot\tilde{p}n\cdot k + \tilde{p}_\perp^2/n\epsilon} \frac{n}{2}$$

soft vertex

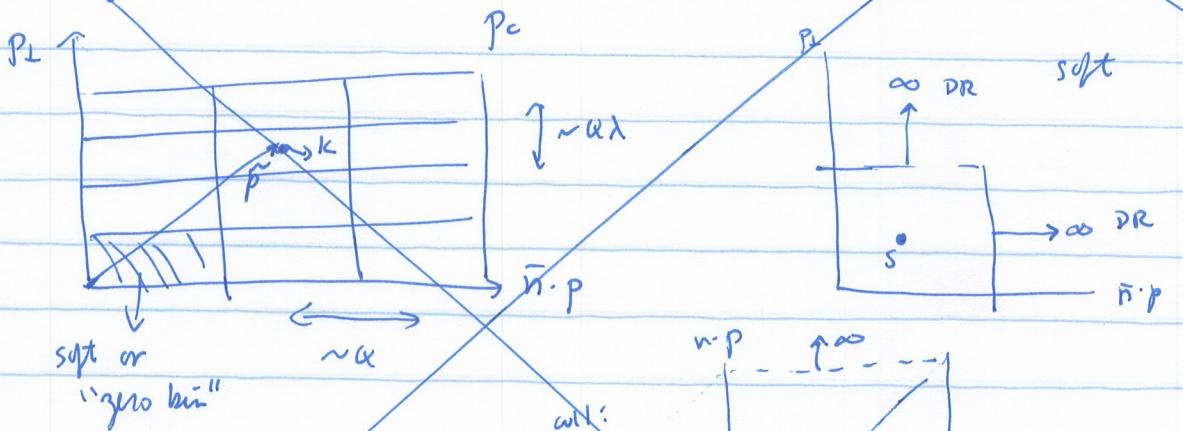
$$\gamma_\mu = \bar{n} \cdot \tilde{p} \frac{n}{2}$$

$$i \frac{\bar{n} \cdot \tilde{p}}{\bar{n} \cdot \tilde{p} n \cdot k n \epsilon} \frac{n}{2} \gamma_\mu \quad i \frac{\bar{n} \cdot \tilde{p}}{n \cdot k n \epsilon} \left(-\gamma_\mu \frac{n}{2} + n_\mu \right) u_n$$

extended prop.

$\sim n^\mu$

picture of mom. space:



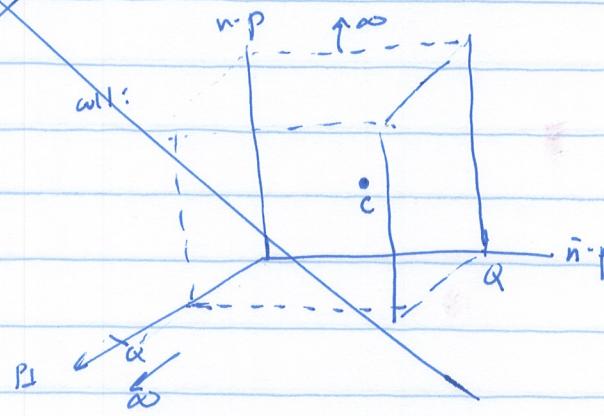
typically,

$$\sum_{\tilde{p} \neq 0} \int d^4k \rightarrow \int d^4p_c (\delta_0 - \delta_{p_0})$$

limit of coll amp

as $p_c \rightarrow 0$.

"kin subscript"



step 1) $\mathcal{L}_{\text{QCD}} = \bar{\psi} i \not{D} \psi + (\text{quarks})$

$$\psi(x) = \sum_{\tilde{p}} e^{-i\tilde{p} \cdot x} \bar{\psi}_{n,\tilde{p}}(x)$$

↳ need sum due to label change: $\frac{\tilde{p}_-}{\tilde{p}-\tilde{p}}$

gluons $A^{\mu}(x) \rightarrow A_c^{\mu}(x) + A_s^{\mu}(x)$

$$A_c^{\mu}(x) = \sum_{\tilde{p}} e^{i\tilde{p} \cdot x} A_{n,\tilde{p}}^{\mu}(x)$$

$$\begin{aligned} \Rightarrow \mathcal{L} &= \sum_{\substack{\tilde{p}, \tilde{p}' \\ \tilde{p} \neq \tilde{p}'}} \bar{\psi}_{n,\tilde{p}} e^{i\tilde{p}' \cdot x} i \not{D} e^{-i\tilde{p} \cdot x} \psi_{n,\tilde{p}} \\ &= \sum_{\tilde{p}, \tilde{p}'} e^{i(\tilde{p}' - \tilde{p}) \cdot x} \bar{\psi}_{n,\tilde{p}} (i \not{D} + \tilde{p}') \psi_{n,\tilde{p}} \end{aligned}$$

$\gamma^0 = \begin{pmatrix} 0 & 0 \\ 0 & 1_0 \end{pmatrix}, \gamma^3 = \begin{pmatrix} 0 & \vec{n} \cdot \vec{\sigma} \\ \vec{n} \cdot \vec{\sigma} & 0 \end{pmatrix}$ 2) define projectors $P_n = \frac{\gamma^0 \gamma^3}{4}, P_{\bar{n}} = \frac{\gamma^3 \gamma^0}{4}$

$$P_n + P_{\bar{n}} = 1$$

$$\frac{\gamma^0 \gamma^3}{4} + \frac{\gamma^3 \gamma^0}{4} = \cancel{\gamma^0 \gamma^0} \frac{1}{4} + 2 \frac{\vec{n} \cdot \vec{n}}{4} = 1$$

note

$$P_n^+ = \frac{\vec{n} + \vec{n}^+}{4} = \frac{\sqrt{m}}{4} = P_n$$

$$P_n^+ \gamma_0 = \gamma_0 \frac{\sqrt{m}}{4}$$

$$\bar{\xi}_n = \bar{\psi}_n \frac{\sqrt{m}}{4} = \bar{\psi}_n P_{\bar{n}}$$

$$\bar{\psi}_{n,\tilde{p}} = \sum_{\substack{\text{III} \\ P_n \bar{\psi}_{n,\tilde{p}}} \bar{\psi}_{n,\tilde{p}}} + \sum_{\substack{\text{III} \\ P_{\bar{n}} \bar{\psi}_{n,\tilde{p}}} \bar{\psi}_{n,\tilde{p}}}$$

$$P_n P_n = \frac{\gamma^0 \gamma^3}{4} \frac{\gamma^0 \gamma^3}{4} = \frac{\gamma^0 \gamma^0}{4} 2 \frac{\vec{n} \cdot \vec{n}}{4} = \frac{\gamma^0 \gamma^0}{4} = P_n$$

$$P_{\bar{n}}^2 = P_{\bar{n}}, P_n P_{\bar{n}} = 0$$

note $\sqrt{m} \bar{\psi}_{n,\tilde{p}} = 0$ (Dirac eq for $\tilde{p} = \frac{\vec{n} \cdot \vec{p}}{2} n^m$)

$$\mathcal{L} = \sum_{\tilde{p}, \tilde{p}'} e^{i(\tilde{p}' - \tilde{p}) \cdot x} (\bar{\xi}_{n,\tilde{p}'} + \bar{\Xi}_{n,\tilde{p}'}) (i \not{D} + \tilde{p}') (\xi_{n,\tilde{p}} + \Xi_{n,\tilde{p}})$$

$$\gamma^M = \sqrt{m} \frac{n^m}{2} + \sqrt{m} \frac{\vec{n} \cdot \vec{\sigma}}{2} + \gamma_L^M \quad \text{use} \quad P_{\bar{n}} \gamma^M P_n = \frac{\sqrt{m}}{4} \gamma^M \frac{\sqrt{m}}{4} = \frac{\sqrt{m}}{4} \sqrt{m} \frac{\sqrt{m}}{4} \frac{n^m}{2} = P_{\bar{n}} \sqrt{m} \frac{n^m}{2}$$

$$P_n \gamma^M P_{\bar{n}} = \frac{\sqrt{m}}{4} \sqrt{\frac{n^m}{2}} \frac{\sqrt{m}}{4} = P_n \frac{\sqrt{m}}{2} \frac{n^m}{2}$$

$$P_n \gamma^M P_n = P_n \gamma_L^M P_n$$

$$P_{\bar{n}} \gamma^M P_{\bar{n}} = P_{\bar{n}} \gamma_L^M P_{\bar{n}}$$

$$\begin{aligned} \Rightarrow \mathcal{L} &= \sum_{\tilde{p}, \tilde{p}'} e^{i(\tilde{p}' - \tilde{p}) \cdot x} \left[\bar{\xi}_{n,\tilde{p}'} \frac{\sqrt{m}}{2} i \not{D} \xi_{n,\tilde{p}} + \bar{\Xi}_{n,\tilde{p}'} \frac{\sqrt{m}}{2} (i \not{D} + \vec{n} \cdot \vec{p}) \Xi_{n,\tilde{p}} \right. \\ &\quad \left. + \bar{\xi}_{n,\tilde{p}'} (i \not{D} + \vec{p}_L) \Xi_{n,\tilde{p}} + \bar{\Xi}_{n,\tilde{p}'} (i \not{D} + \vec{p}_L) \xi_{n,\tilde{p}} \right] \end{aligned}$$

3) $E_{n,p}$ is a heavy field w/ "mass" $\bar{n} \cdot \tilde{p} \sim \Omega$
 int. mt. : eq. of motion

$$(i\bar{n} \cdot D + \bar{n} \cdot \tilde{p}) \frac{\not{k}}{2} \xi_{n,p} = - (iD_L + \tilde{p}_L) \xi_{n,p}$$

$$\Rightarrow \boxed{\xi_{n,p} = \frac{iD_L + \tilde{p}_L}{i\bar{n} \cdot D + \bar{n} \cdot \tilde{p}} \frac{\not{k}}{2} \xi_{n,p}}$$

$$\Rightarrow \mathcal{L} = \sum_{p,p'} e^{i(\tilde{p}-\tilde{p}') \cdot x} \bar{\xi}_{n,p'} [in \cdot D + (iD_L + \tilde{p}_L) \frac{1}{i\bar{n} \cdot D + \bar{n} \cdot \tilde{p}} (iD_L + \tilde{p}_L)] \frac{\not{k}}{2} \xi_{n,p}$$

4) expand in λ : note $A_c^m \sim \alpha(-, \lambda^2, \lambda)$ just like their momenta
 $A_s^m \sim \alpha \Omega^2, \lambda^2, \lambda^2$

$$\Rightarrow in \cdot D = in \cdot \partial + gn \cdot A_c + gn \cdot A_s \sim \alpha \lambda^2$$

$$iD_L + \tilde{p}_L = i\tilde{p}_L + gn \cdot A_L^c + gn \cdot A_L^s \sim \alpha \lambda$$

$$i\bar{n} \cdot D + \bar{n} \cdot \tilde{p} = i\bar{n} \cdot \partial + g\bar{n} \cdot A_c + g\bar{n} \cdot A_s \sim \alpha$$

define "local operator" $P^\mu d_{n,p} = \tilde{p}^\mu d_{n,p}$

$$P^\mu d_{n,p}^* = -\tilde{p}^\mu d_{n,p}^*$$

$$\text{and } iD_c^m \equiv P^\mu + gA_c^m$$

$$\Rightarrow \mathcal{L}_{\text{coll}} = e^{iP \cdot x} \bar{\xi}_{n,p} [in \cdot D + iD_c^m \frac{1}{i\bar{n} \cdot D_c} iD_c^m] \frac{\not{k}}{2} \xi_{n,p}$$

gives coll. prop
& coll-soft vertex $--\frac{\xi}{n_p}--$

5) can rewrite last term w/ Wilson lines:

$$W_n \stackrel{(x)}{=} P \exp \left(ig \int_{-\infty}^x ds \bar{n} \cdot A_n (\bar{n}s) \right)$$

$$\text{obey} \quad i\bar{n} \cdot D_c W_n = 0$$

$$\begin{aligned} i\bar{n} \cdot D_c (W_n \phi) &= (\bar{\rho} + g n \cdot A_n) (W_n \phi) \\ &= [i\bar{n} \cdot D_c W_n] \phi + W_n \bar{\rho} \phi \\ &\stackrel{!}{=} 0 \end{aligned}$$

$$\Rightarrow i\bar{n} \cdot D_n (W_n \phi) = W_n \bar{\rho} \phi$$

$$\Rightarrow W_n^+ i\bar{n} \cdot D_n W_n = \bar{\rho}$$

$$i\bar{n} \cdot D_n = W \bar{\rho} W^+$$

$$\Rightarrow \tilde{\chi} = e^{-i\bar{\rho} \cdot x} \sum_{pp'} \tilde{f}_{n,p'} (i\bar{n} \cdot D + iB_L^c W \frac{1}{\bar{\rho}} W^+ iB_L^c) \frac{\sqrt{\pi}}{2} \delta_{n,p}$$

χ
coll soft
decaying ...
next lecture