

01 Jun 2016

Lecture 4: Factorization in SCET_I

(see lecture notes online)

remainder of lectures: * 4) factorization, Wilson lines, soft NP universality

* { 5) 1-loop matching calc.
6) 1-loop soft func. calc., RG & resummation

⚡ { 7) DIS, pp & beam functions, NP universality
8) SCET_{II}, 3, + ...

Yesterday ended with

$$\mathcal{L} = e^{iP \cdot x} \int \bar{\psi}_n [i\bar{n} \cdot D + i\cancel{D}_\perp^c \frac{1}{i\bar{n} \cdot D_c} i\cancel{D}_\perp^c] \psi_n$$

$$\bar{\psi}_n = \int_{\bar{p}} \bar{\psi}_n \tilde{p} \quad A_n = \int_{\tilde{p}} A_n \tilde{p}$$

$$i\bar{n} \cdot D = i\bar{n} \cdot \partial + g\bar{n} \cdot A_n + g\bar{n} \cdot A_s$$

$$i\cancel{D}_\perp^{\mu} = P^\mu + g A_n^\mu$$

rewrite last term: use coll. Wilson line

$$W_n(x) = \underset{\text{path ordering}}{P} \exp \left[ig \int_{-\infty}^x ds \bar{n} \cdot A_n(\bar{n}s) \right]$$

$$i\bar{n} \cdot D_c W_n = 0$$

$$\text{so } i\bar{n} \cdot D_c (W_n \psi_c) = \cancel{(i\bar{n} \cdot D_c W_n)} \psi_c + W_n \bar{P} \psi_c$$

$$\Rightarrow [i\bar{n} \cdot D_c W_n \cdot] \stackrel{\text{operator}}{=} W_n \bar{P}$$

$$\Rightarrow \bar{n} \cdot D_c = W_n \bar{P} W_n^\dagger$$

$$\Rightarrow \frac{1}{i\bar{n} \cdot D_c} = W_n \frac{1}{\bar{P}} W_n^\dagger$$

$$\Rightarrow \mathcal{L} = e^{iP \cdot x} \int \bar{\psi}_n [i\bar{n} \cdot D + i\cancel{D}_\perp^c W_n \frac{1}{\bar{P}} W_n^\dagger i\cancel{D}_\perp^c] \psi_n$$

Feynman rules: $-\cdots\rightarrow-\cdots = i \frac{\kappa}{2} \frac{\bar{n} \cdot p}{\bar{n} \cdot p \cdot k + p_1^2 \cdot n_1 \cdot \epsilon}$

$-\cdots\left\{\begin{matrix} \text{CS} \\ \text{---} \end{matrix}\right. = i g t^a n_\mu \frac{\kappa}{2}$

($e^{i p \cdot x}$ enters label man. Conservation at vertices)

$-\rightarrow\left\{\begin{matrix} \text{---} \\ \text{---} \end{matrix}\right.\rightarrow = i g t^a \left[n_\mu + \frac{\delta_\mu^+ \beta_1}{\bar{n} \cdot p} + \frac{\beta_1' \delta_\mu^+}{\bar{n} \cdot p'} - \frac{\beta_1' \beta_1}{\bar{n} \cdot p \cdot \bar{n} \cdot p'} \bar{n}_\mu \right] \frac{\kappa}{2}$

$\begin{matrix} \mu, a & \nu, b \\ \text{---} & \text{---} \\ \text{---} & \text{---} \end{matrix} = \frac{g^2 t^a t^b}{\bar{n} \cdot (p-q)} \left[\gamma_\mu^+ \gamma_\nu^+ - \frac{\gamma_\mu^+ \beta_1}{\bar{n} \cdot p} \bar{n}_\nu - \frac{\beta_1' \gamma_\nu}{\bar{n} \cdot p'} \bar{n}_\mu + \frac{\beta_1' \beta_1}{\bar{n} \cdot p \cdot \bar{n} \cdot p'} \bar{n}_\mu \bar{n}_\nu \right] \frac{\kappa}{2}$

int. w/ $\Xi_n!$ $+ \frac{1}{\bar{n} \cdot (q,p)} (a \leftrightarrow b, \mu \leftrightarrow \nu)$

soft-collinear decoupling:

look at $\bar{\xi}_n (in \cdot D_S) \xi_n$ $i D_S = i \partial^\mu + g A_S^\mu$

can define soft Wilson line $Y_n(x) = P \exp \left[i g \int_0^x ds n \cdot A_S(s) \right]$
 $in \cdot D_S Y_n = 0$

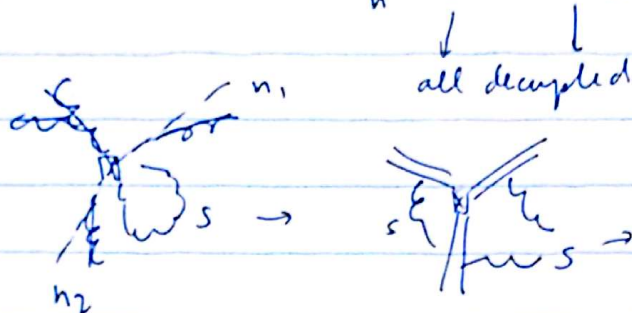
field redefinition: $\bar{\xi}_n(x) = Y_n(x) \bar{\xi}_n^{(0)}(x)$ $A_n(x) = Y_n A_n^{(0)} Y_n^\dagger$
 $\Rightarrow W_n = Y_n W_n^{(0)} Y_n^\dagger$

$\Rightarrow \bar{\xi}_n in \cdot D_S \xi_n \rightarrow \bar{\xi}_n^{(0)} Y_n^\dagger in \cdot D_S Y_n \xi_n^{(0)}$
 $= \bar{\xi}_n^{(0)} Y_n^\dagger Y_n in \cdot \partial \frac{\kappa}{2} \xi_n^{(0)}$
 $= \bar{\xi}_n^{(0)} in \cdot \partial \frac{\kappa}{2} \xi_n^{(0)}$
 no soft gluons!

similarly $\bar{\xi}_n g \bar{n} \cdot A_n \xi_n \rightarrow \bar{\xi}_n^{(0)} Y_n^\dagger Y_n \bar{n} \cdot A_n^{(0)} Y_n^\dagger Y_n \xi_n^{(0)} = \bar{\xi}_n^{(0)} \bar{n} \cdot A_n^{(0)} \xi_n^{(0)}$
 etc.

similarly in coll. gluon leg.

This means $\mathcal{L}_{scet} = \sum_n \mathcal{L}_{c,n}^{(0)} + \mathcal{L}_S$



all decoupled
cannot resolve collinear splittings
only total n_i & color
(3 or 8)

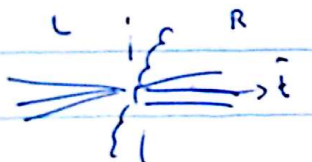
in fact $\mathcal{L}_{c,n}^{(0)}$ is just $\mathcal{L}_{c,0}$ boosted in n direction

so \mathcal{L}_{scet} is sum of decoupled copies of dCD in each coll.
& soft sector, but at different scales.

2-jet factorization

Consider a measurement of "thrust" in e^+e^- :

$$e^+e^- \rightarrow X$$

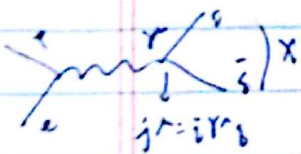


$$T = \frac{1}{Q} \max_{\hat{e}} \sum_{i \in X} |\hat{e} \cdot \vec{p}_i| = \frac{1}{Q} (\vec{P}_z^L + \vec{P}_z^R)$$

$\rightarrow 1$ for back-to-back

$$\text{use } \tau = 1 - T$$

$\rightarrow 0$



$$\frac{d\Gamma}{d\tau} = \frac{1}{2Q^2} \sum_X |\langle X | \bar{\psi} \gamma^\mu \psi | 0 \rangle L_\mu|^2 (2\pi)^4 \delta^4(Q - p_X) \delta(\tau - \tau(X))$$

L_μ = leptonic tensor

operator matching $j^\mu = \bar{\psi} \gamma^\mu \psi \rightarrow j_{scet}^\mu = \sum_{n_1, n_2} C_2(\vec{P}_1, \vec{P}_2, \mu) \bar{\chi}_{n_1} \gamma^\mu \chi_{n_2}$

$$\chi_{n_i} = W_{n_i}^\dagger \xi_{n_i}$$

actually

after field redef.

$$C_2(\vec{P}_1, \vec{P}_2, \mu) = \bar{\chi}_{n_2} \gamma^\mu \chi_{n_1}$$

$$j^\mu \rightarrow C_2(p) O_2(p) = \bar{\chi}_{n_2} \gamma^\mu \chi_{n_1} C(\vec{P}_{n_1}, \vec{P}_{n_2}, \mu) \chi_{n_1} \chi_{n_2}$$

$$\sum_{\substack{Y_n, Y_s \\ X_s}} |X_n Y_n X_s\rangle \langle Y_n Y_n X_s|$$

$$\Rightarrow \frac{d\Gamma}{d\Omega} = \frac{1}{2\omega^2} L_\mu L_\nu \sum_X \langle 0 | \bar{X}_{n_1} Y_{n_1}^\dagger \mathcal{E}(P_1, P_2, P) \hat{Y}_{n_2} X_{n_2} | X \rangle$$

$$\times \langle X | \bar{X}_{n_2} Y_{n_2}^\dagger \mathcal{E}(P_1, P_2, P) \hat{Y}_{n_1} X_{n_1} | 0 \rangle$$

$$\times (2\pi)^4 \delta^4(Q - P_X) \delta(\tau - 1 + \frac{|P_2^L + P_2^R|}{Q})$$

n_1, n_2 back-to-back

$$(Q, \vec{0}) = P_{n_1} + P_{n_2} + P_S$$

$$(Q, \tau Q, \vec{0}_\perp) = (P_1^-, P_1^+, P_1^\perp) \Rightarrow \tilde{P}_1^- = Q$$

$$+ (P_2^-, P_2^+, P_2^\perp) \Rightarrow \tilde{P}_2^+ = Q$$

$$+ (P_S^+, P_S^-, P_S^\perp) \Rightarrow \tilde{P}_1^\perp = -\tilde{P}_2^\perp \rightarrow 0$$

residual & soft moment unstrained. ($k \rightarrow \infty$)
(by $\delta^4(Q - p_i)$, anyway)

~~measurement: $\tau = 1 - \frac{P_2^L + P_2^R}{Q}$~~

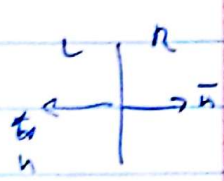
~~$Q\tau = Q - P_2^L - P_2^R$~~

~~$= Q - (P_1^L + P_2^L + P_S^L) - (P_{c1}^R + P_{c2}^R + P_S^R)$~~

~~$= Q - (\frac{P_1^- - P_1^+}{2}) - (\frac{P_2^+ - P_2^-}{2})$~~

~~$= Q - \tilde{P}_1^-$~~

~~$P_2 = \frac{P_2^+ - P_2^-}{2}$~~



split measurement into coll & soft parts:

$$\tau = \tau_n + \tau_{\bar{n}} + \tau_S$$

Integrating $\delta(\tau - \tau_n - \tau_{\bar{n}} - \tau_S)$

$$\frac{d\Gamma}{d\Omega} = \frac{1}{2\omega^2} L_\mu L_\nu \sum_{\substack{Y_n, Y_s \\ X_n}} \langle 0 | \bar{X}_{n_1}^S | \bar{n} \rangle \langle \bar{n} | X_{n_1}^A | 0 \rangle \delta(\tau_{\bar{n}} - \tau(X_{\bar{n}}))$$

$$\times \sum_{X_n} \langle 0 | Y_{n_1}^+ | X_n \rangle \langle X_n | \bar{X}_{n_1}^e | 0 \rangle \delta(\tau_n - \tau(X_n))$$

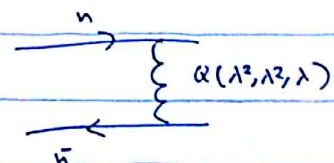
$$\times \sum_{Y_s} \langle 0 | Y_{n_1}^+ Y_{n_2}^- | X_s \rangle \langle X_s | Y_{n_1}^+ Y_{n_2}^- | 0 \rangle \delta(\tau_S - \tau(X_s))$$

$$\times \mathcal{E}^4(Q, P) \mathcal{E}(Q, P)$$

$$\begin{aligned} \tau(\chi_n) &\approx \alpha \tau = Q - \frac{P_L^+ - P_L^-}{2} - \frac{P_R^- - P_R^+}{2} \\ &= P_L^- + P_R^+ \quad (\text{small comp. } v_2 + s/t) \end{aligned}$$

$$\begin{aligned} \Rightarrow \frac{d\sigma}{d\tau} &= \frac{1}{4\alpha^2} \int d\tau_n d\tau_s \delta(\tau - \tau_n - \tau_s) \\ &\times \sum_{\bar{\chi}_n} \langle 0 | \bar{\chi}_n^\delta | \chi_n \rangle \langle \chi_n | \chi_{n,1}^\delta | 0 \rangle \delta(\tau_n - \frac{P_{\chi_n^+}}{\alpha}) \\ &\times \sum_{\chi_n} \langle 0 | \chi_n^\delta | \chi_n \rangle \langle \chi_n | \bar{\chi}_n^\delta | 0 \rangle \delta(\tau_n - \frac{P_{\chi_n^+}}{\alpha}) \\ &\times \sum_{\chi_s} \langle 0 | \chi_n^\dagger \chi_n | \chi_s \rangle \langle \chi_s | \chi_n^\dagger \chi_n | 0 \rangle \delta(\tau_s - \frac{k_s^+}{\alpha} - \frac{k_s^-}{\alpha}) \\ &\times |C|^2 \end{aligned}$$

$$\begin{aligned} \rightarrow \sigma_0 \cdot H(\alpha^2, \mu) \int d\tau_n d\tau_s \delta(\tau - \tau_n - \tau_s) \\ \times J_n(\tau_n) J_n(\tau_s) S(\tau_s, \mu) \\ \downarrow \\ \text{1-loop calc. \& RGE tomorrow.} \end{aligned}$$

sublity : Numbers  off-shell
can include fact.
can include in SCET.

$$\frac{1}{N_c} \text{Tr} \sum_{\chi_n} \langle 0 | \chi_{n,1}^\delta | \chi_n \rangle \langle \chi_n | \bar{\chi}_n^\delta | 0 \rangle \delta(\tau_n - \frac{P_{\chi_n^+}}{\alpha}) = J(\alpha^2, \mu) \left(\frac{\mu}{2}\right)^{P^E}$$

sim. for \bar{n} \downarrow
PP⁺

$$\frac{1}{N_c} \text{Tr} \sum_{\chi_s} \langle 0 | \chi_n^\dagger \chi_n | \chi_s \rangle \langle \chi_s | \chi_n^\dagger \chi_n | 0 \rangle \delta(\tau_s - \frac{k_s^+}{\alpha} - \frac{k_s^-}{\alpha}) = S(k_s^+ + k_s^-)$$

\downarrow
 $k_L^+ + k_n^+$

$$\begin{aligned} \frac{H}{\alpha^2} & \quad P^2 \\ \frac{J}{\alpha^2 \tau} & \\ \frac{S}{\alpha^2 \tau^2} & \end{aligned}$$